

# III-harmonic Curves in $\widetilde{\mathbf{SL}_2\mathbb{R}}$ Space

Bendehiba Senoussi

**Abstract** Some work has been done in the study of non-geodesic III-harmonic curves in some model spaces. In this paper, we study III-harmonic curves in  $\widetilde{\mathbf{SL}_2\mathbb{R}}$  space. We give necessary and sufficient conditions for helices to be III-harmonic. Also, we characterize III-harmonic curves in terms of their curvature and torsion.

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## 1 Introduction

First we should recall some notions and results related to the harmonic and the polyharmonic ( $r$ -harmonic  $r \geq 1$ ) maps between Riemannian manifolds. Harmonic maps  $\psi : (\mathcal{M}, g) \rightarrow (\mathcal{N}, \tilde{g})$  between Riemannian manifolds are the critical points of the energy functional

$$\mathcal{E}_1 : C^\infty(\mathcal{M}, \mathcal{N}) \rightarrow \mathbb{R}, \quad \mathcal{E}_1(\psi) = \frac{1}{2} \int_{\mathcal{M}} |\mathrm{d}\psi|^2 v_g,$$

and are characterized by the vanishing of the first tension field

$$\tau_1(\psi) = -\mathrm{d}^*\mathrm{d}\psi = \operatorname{trace} \nabla \mathrm{d}\psi,$$

where  $\mathrm{d}$  is the exterior differentiation and  $\mathrm{d}^*$  is the codifferentiation.

We remind that the bienergy of  $\psi$  is given by

$$\mathcal{E}_2 : C^\infty(\mathcal{M}, \mathcal{N}) \rightarrow \mathbb{R}, \quad \mathcal{E}_2(\psi) = \frac{1}{2} \int_{\mathcal{M}} |\tau(\psi)|^2 v_g,$$

and the bitension field  $\tau_2(\psi)$  has the expression

$$\tau_2(\psi) = -\Delta^\psi \tau(\psi) - \operatorname{trace}_g \mathbf{R}^{\mathcal{N}}(\mathrm{d}\psi, \tau(\psi))\mathrm{d}\psi,$$

where  $\Delta^\psi = -\text{trace}(\nabla^\psi)^2 = -\text{trace}(\nabla^\psi \nabla^\psi - \nabla_{\nabla^\psi}^\psi)$ .

A smooth map  $\psi$  is *biharmonic* if it satisfies the following biharmonic equation

$$\tau_2(\psi) = 0.$$

Biharmonic maps are the critical points of the bienergy functional  $\mathcal{E}_2$ . We call proper biharmonic the non-harmonic biharmonic maps. Biharmonic curves  $\psi$  of a Riemannian manifold are the solutions of the fourth order differential equation

$$\nabla_{\psi'}^3 \psi' - R(\psi', \nabla_{\psi'} \psi') \psi' = 0. \quad (1.1)$$

Eells and Lemaire [8] proposed the problem to consider the polyharmonic ( $r$ -harmonic  $r \geq 1$ ) maps of order  $r$ , these are critical points of the  $r$ -energy functional defined by

$$\mathcal{E}_r(\psi) = \int_{\mathcal{M}} e_r(\psi) v_g, \quad r \geq 1, \quad (1.2)$$

where  $e_r(\psi) = \frac{1}{2} \|(d + d^*)^r \psi\|^2$  for smooth maps  $\psi$ . A map  $\psi$  is  $r$ -harmonic if it is a critical point of the functional  $\mathcal{E}_r(\psi)$  defined in (1.2).

Every harmonic map is a solution of the polyharmonic map, see [4] for a recent classification result. In [27], S. B. Wang studied the first variation formula of the  $k$ -energy  $\mathcal{E}_k$  whose critical maps are called  $k$ -harmonic maps. In [13], S. Maeta showed the second variation formula of the  $k$ -energy. Triharmonic curves with constant curvature in space forms were studied by Maeta in [13]. Recently, for the study of triharmonic curves in 3-dimensional homogeneous spaces see [21–23] respectively. The main goal of this article is to define and characterize III-harmonic curves in  $\widetilde{\mathbf{SL}_2\mathbb{R}}$  space.

## 2 Preliminaries

In [14] Emil Molnár proposed a projective spherical model as unified geometrical model of homogeneous geometries.  $\widetilde{\mathbf{SL}_2\mathbb{R}}$  geometry is one of the eight homogeneous Thurston 3-geometries. The hyperboloid model of  $\widetilde{\mathbf{SL}_2\mathbb{R}}$  geometry is described in details in [7, 10, 14–16]. In  $\widetilde{\mathbf{SL}_2\mathbb{R}}$ , the metric is given by

$$ds^2 = dr^2 + (\cosh r \sinh r)^2 d\vartheta^2 + (d\phi + \sinh^2 r d\vartheta)^2,$$

where  $\phi$  is a fiber coordinate and  $(r, \vartheta)$  are polar coordinates of the intersection point of a fiber and the hyperbolic base plane.

If

$$\{\theta^1 = dr, \theta^2 = \cosh r \sinh r d\vartheta, \theta^3 = d\phi + \sinh^2 r d\vartheta\}$$

is an orthonormal coframe then the dual orthonormal frame fields are given by

$$\begin{pmatrix} e_1 \\ e_2 \\ e_3 \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \frac{2}{\sinh 2r} & -\tanh r \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \partial_r \\ \partial_{\vartheta} \\ \partial_{\phi} \end{pmatrix},$$

where  $\partial_r = \frac{\partial}{\partial r}$ ,  $\partial_\vartheta = \frac{\partial}{\partial \vartheta}$ ,  $\partial_\phi = \frac{\partial}{\partial \phi}$ .

One can easily check that  $g(e_i, e_j) = \delta_{ij}$  where  $\delta_{ij}$  denotes the Kronecker delta.

The symmetric metric tensor field  $g$  is given by

$$G = (g_{ij})_{1 \leq i, j \leq 3} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \sinh^2 r (\cosh^2 r + \sinh^2 r) & \sinh^2 r \\ 0 & \sinh^2 r & 1 \end{pmatrix}.$$

The inverse matrix  $(g^{ij})_{1 \leq i, j \leq 3}$  of  $(g_{ij})_{1 \leq i, j \leq 3}$  is given by

$$G^{-1} = \frac{4}{\sinh^2 2r} \begin{pmatrix} \sinh^2 r \cosh^2 r & 0 & 0 \\ 0 & 1 & -\sinh^2 r \\ 0 & -\sinh^2 r & \sinh^2 r (\cosh^2 r + \sinh^2 r) \end{pmatrix}.$$

The Levi-Civita connection  $\nabla$  is defined by  $\nabla_{\partial_i} \partial_j = \Gamma_{ij}^k \partial_k$ , where the Cristoffel symbols  $\Gamma_{ij}^k$  are given by

$$\begin{aligned} \Gamma_{ij}^1 &= -\frac{1}{2} \sinh 2r \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 + 4 \sinh^2 r & 1 \\ 0 & 1 & 0 \end{pmatrix}, \\ \Gamma_{ij}^2 &= \frac{2}{\sinh 2r} \begin{pmatrix} 0 & 1 + 3 \sinh^2 r & 1 \\ 1 + 3 \sinh^2 r & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix}, \\ \Gamma_{ij}^3 &= -\tanh r \begin{pmatrix} 0 & 2 \sinh^2 r & 1 \\ 2 \sinh^2 r & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix}. \end{aligned}$$

The Levi-Civita connection  $\nabla$  is given by

$$\begin{aligned} \nabla_{e_1} \begin{pmatrix} e_1 \\ e_2 \\ e_3 \end{pmatrix} &= \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} e_1 \\ e_2 \\ e_3 \end{pmatrix}, \\ \nabla_{e_2} \begin{pmatrix} e_1 \\ e_2 \\ e_3 \end{pmatrix} &= \begin{pmatrix} 0 & 2 \coth 2r & 1 \\ -2 \coth 2r & 0 & 0 \\ -1 & 0 & 0 \end{pmatrix} \begin{pmatrix} e_1 \\ e_2 \\ e_3 \end{pmatrix}, \\ \nabla_{e_3} \begin{pmatrix} e_1 \\ e_2 \\ e_3 \end{pmatrix} &= \begin{pmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} e_1 \\ e_2 \\ e_3 \end{pmatrix}. \end{aligned} \tag{2.1}$$

The characterising properties of this algebra are the following commutation relations:

$$[e_1, e_2] = -2 \coth(2r) e_2 - 2e_3, \quad [e_1, e_3] = [e_2, e_3] = 0.$$

The Riemannian curvature tensor is given by

$$R(X_1, X_2, X_3, X_4) = g_{\widetilde{\mathbf{SL}_2 \mathbb{R}}} (R(X_2, X_1)X_3, X_4) = -g_{\widetilde{\mathbf{SL}_2 \mathbb{R}}} (R(X_1, X_2)X_3, X_4), \tag{2.2}$$

where  $X_1, X_2, X_3, X_4$  are smooth vector fields on  $\widetilde{\mathbf{SL}_2\mathbb{R}}$ .

The components of the Riemannian curvature tensor defined by

$$R(X_1, X_2)X_3 = (\nabla_{X_1}\nabla_{X_2} - \nabla_{X_2}\nabla_{X_1} - \nabla_{[X_1, X_2]})X_3$$

are [9]

$$\begin{aligned} R(e_1, e_2)e_1 &= 7e_2, & R(e_1, e_2)e_2 &= -7e_1, & R(e_1, e_3)e_1 &= -e_3, \\ R(e_1, e_3)e_3 &= e_1, & R(e_3, e_2)e_2 &= e_3, & R(e_3, e_2)e_3 &= -e_2. \end{aligned} \quad (2.3)$$

The Riemannian curvature tensor field  $R$  is described by

$$R_{1212} = -7, \quad R_{1313} = R_{2323} = 1. \quad (2.4)$$

The mixed product of three vectors  $X, Y, Z$  of a  $\widetilde{\mathbf{SL}_2\mathbb{R}}$  is

$$[X, Y, Z] = g_{\widetilde{\mathbf{SL}_2\mathbb{R}}}(X \wedge Y, Z). \quad (2.5)$$

### 3 Polyharmonic curves in $\widetilde{\mathbf{SL}_2\mathbb{R}}$ space

#### 3.1 Biharmonic curves in $\widetilde{\mathbf{SL}_2\mathbb{R}}$ space

In [6] and [11], the authors studied biharmonic curves in BCV-spaces and they obtained interesting classification results. In [5] the authors proved that any proper biharmonic curve in the Heisenberg group  $\mathcal{H}_3$  is a helix and gave its explicit parametrization.

Let  $\psi : I \rightarrow \widetilde{\mathbf{SL}_2\mathbb{R}}$  be a differentiable curve parametrized by arc length and let  $\{T = \psi'(s), N, B\}$  be the orthonormal moving Frenet frame along the curve  $\psi$  in  $\widetilde{\mathbf{SL}_2\mathbb{R}}$  such that  $T = \psi'$  is the unit vector field tangent to  $\psi$ ,  $N$  is the unit vector field in the direction  $\nabla_T T$  normal to  $\psi$  (principal normal) and  $B = T \wedge N$  (binormal vector). Then we have the following Frenet equations:

$$\nabla_T \begin{pmatrix} T \\ N \\ B \end{pmatrix} = \begin{pmatrix} 0 & k & 0 \\ -k & 0 & \tau \\ 0 & -\tau & 0 \end{pmatrix} \begin{pmatrix} T \\ N \\ B \end{pmatrix}, \quad (3.1)$$

where

$$\begin{aligned} g_{\widetilde{\mathbf{SL}_2\mathbb{R}}}(T, T) &= g_{\widetilde{\mathbf{SL}_2\mathbb{R}}}(N, N) = g_{\widetilde{\mathbf{SL}_2\mathbb{R}}}(B, B) = 1, \\ g_{\widetilde{\mathbf{SL}_2\mathbb{R}}}(T, N) &= g_{\widetilde{\mathbf{SL}_2\mathbb{R}}}(T, B) = g_{\widetilde{\mathbf{SL}_2\mathbb{R}}}(N, B) = 0, \end{aligned}$$

and  $k$  ( $k^2 = g_{\widetilde{\mathbf{SL}_2\mathbb{R}}}(\nabla_T T, \nabla_T T)$ ) is the curvature of  $\psi$  and  $\tau$  is its torsion.

From (3.1) we have

$$\nabla_T^3 T = (-3kk')T + (k'' - k^3 - k\tau^2)N + (2k'\tau + k\tau')B, \quad (3.2)$$

where  $k' = \frac{dk}{ds}$ ,  $k'' = \frac{d^2k}{ds^2}$ ,  $\tau' = \frac{d\tau}{ds}$ .

With respect to the orthonormal basis  $\{e_1, e_2, e_3\}$  we can write

$$\begin{pmatrix} T \\ N \\ B \end{pmatrix} = \begin{pmatrix} T_1 & T_2 & T_3 \\ N_1 & N_2 & N_3 \\ B_1 & B_2 & B_3 \end{pmatrix} \begin{pmatrix} e_1 \\ e_2 \\ e_3 \end{pmatrix}. \quad (3.3)$$

A direct computation using (2.4) and (3.3) gives the following:

$$\begin{cases} R(T, N, T, N) = \sum_{1 \leq i, j, k, l \leq 3} T_i N_j T_k N_l R_{ijkl} = 1 - 8B_3^2 \\ R(T, N, T, B) = \sum_{1 \leq i, j, k, l \leq 3} T_i N_j T_k B_l R_{ijkl} = 8N_3 B_3 \\ R(B, N, T, N) = \sum_{1 \leq i, j, k, l \leq 3} B_i N_j T_k N_l R_{ijkl} = 8T_3 B_3 \\ R(B, N, T, B) = \sum_{1 \leq i, j, k, l \leq 3} B_i N_j T_k B_l R_{ijkl} = -8T_3 N_3 \\ R(B, T, T, B) = \sum_{1 \leq i, j, k, l \leq 3} B_i T_j T_k B_l R_{ijkl} = 8N_3^2 - 1. \end{cases} \quad (3.4)$$

We can use (2.1) and compute the covariant derivatives of the vector fields  $T$ ,  $N$  and  $B$  as:

$$\begin{aligned} \nabla_T T &= (T'_1 - 2T_2^2 \coth 2r - 2T_2 T_3) e_1 + (T'_2 + 2T_1 T_2 \coth 2r + 2T_1 T_3) e_2 + (T'_3) e_3, \\ \nabla_T N &= (N'_1 - 2T_2 N_2 \coth 2r - T_2 N_3 - T_3 N_2) e_1 + (N'_2 + 2T_2 N_1 \coth 2r + T_1 N_3 + T_3 N_1) e_2 \\ &\quad + (N'_3 + T_2 N_1 - T_1 N_2) e_3, \\ \nabla_T B &= (B'_1 - 2T_2 B_2 \coth 2r - T_2 B_3 - T_3 B_2) e_1 + (B'_2 + 2T_2 B_1 \coth 2r + T_1 B_3 + T_3 B_1) e_2 \\ &\quad + (B'_3 + T_2 B_1 - T_1 B_2) e_3. \end{aligned} \quad (3.5)$$

**Theorem 3.1** ([9]). *Let  $\psi : I \rightarrow \widetilde{\mathbf{SL}_2\mathbb{R}}$  be a differentiable curve parametrized by arc length. Then  $\psi$  is a proper non-geodesic biharmonic curve if and only if*

$$\begin{cases} k = \text{constant} \neq 0 \\ k^2 + \tau^2 = 1 - 8B_3^2 \\ \tau' = -8N_3 B_3 \end{cases}. \quad (3.6)$$

### 3.2 III-harmonic curves in $\widetilde{\mathbf{SL}_2\mathbb{R}}$ space

To study III-harmonic curves in  $\widetilde{\mathbf{SL}_2\mathbb{R}}$ , we shall use their Frenet vector fields and equations.

Let us denote by  $\psi : I \rightarrow \mathcal{M}^n$  an arc length parametrized curve in  $\mathcal{M}^n$ . Assume that  $\psi$  is non-geodesic.

If  $r = 2t$ ,  $t \geq 1$ , then (1.2) takes the form [12, 27]

$$\mathcal{E}_{2t}(\psi) = \frac{1}{2} \int_I \left\langle \underbrace{(d^*d) \dots (d^*d)}_{t \text{ times}} \psi, \underbrace{(d^*d) \dots (d^*d)}_{t \text{ times}} \psi \right\rangle v_g. \quad (3.7)$$

If  $r = 2t + 1$ , then (1.2) takes the form

$$\mathcal{E}_{2t+1}(\psi) = \frac{1}{2} \int_I \left\langle d \underbrace{(d^*d) \dots (d^*d)}_{t \text{ times}} \psi, \underbrace{(d^*d) \dots (d^*d)}_{t \text{ times}} \psi \right\rangle v_g. \quad (3.8)$$

The Euler-Lagrange equations of (3.7) and (3.8), reduce to the equation

$$\tau_r(\psi) = \nabla_T^{2r-1}T + \sum_{j=0}^{r-1} (-1)^j R(\nabla_T^{2r-3-j}T, \nabla_T^jT)T, \quad r \geq 1. \quad (3.9)$$

Solutions of equation (3.9) are called as  $r$ -harmonic curves.

The  $r$ -energy functionals defined in (3.7), (3.8) have been intensively studied (see [12, 17–19, 27]).

**Remark 3.1.** i) If  $r = 1$ , the functional (3.8) is just the energy.

ii) If  $r = 2$ , the functional (3.7) is called bienergy and its critical points are the so-called biharmonic maps.

iii)  $r$ -harmonic curves called as proper  $r$ -harmonic if they are not harmonic.

iv) Any harmonic curve is a  $r$ -harmonic curve, for any  $r \geq 1$ .

An arc-length parametrized curve  $\psi : I \rightarrow \mathcal{M}^n$  from  $I \subset \mathbb{R}$  to a Riemannian manifold  $\mathcal{M}^n$  of dimension  $n$  is called *triharmonic* ( $r = 3$ ) if [17]

$$\nabla_T^5T + R(\nabla_T^3T, T)T - R(\nabla_T^2T, \nabla_TT)T = 0. \quad (3.10)$$

With respect to the Frenet frame, the III-harmonic curve equation takes the following form.

**Proposition 3.2.** Let  $\psi : I \subset \mathbb{R} \rightarrow \mathcal{M}^3$  be a differentiable curve parametrized by arc length. Then  $\psi$  is III-harmonic curve if and only if

$$\begin{cases} \zeta_1(s) = 0 \\ \zeta_2(s) - \zeta_4(s)R(N, T, T, N) - \zeta_5(s)R(B, T, T, N) + \zeta_6(s)R(B, N, T, N) = 0 \\ \zeta_3(s) - \zeta_4(s)R(N, T, T, B) - \zeta_5(s)R(B, T, T, B) + \zeta_6(s)R(B, N, T, B) = 0, \end{cases} \quad (3.11)$$

where

$$\begin{aligned} \zeta_1(s) &= -10k'k'' - 5kk^{(3)} + 5kk'(2k^2 + \tau^2) + 5k^2\tau\tau' \\ \zeta_2(s) &= k^5 + k^{(4)} - 15kk'^2 - 10k^2k'' + k\tau^2(2k^2 + \tau^2) - 6\tau^2k'' \\ &\quad - 12k'\tau\tau' - 3k\tau'^2 - 4k\tau\tau'' \\ \zeta_3(s) &= 4\tau k^{(3)} + k\tau^{(3)} - 9k^2k'\tau - 4k'(\tau^3 - \tau'') + 6\tau'(k'' - k\tau^2) - \tau'k^3 \\ \zeta_4(s) &= k'' - k(2k^2 + \tau^2) \\ \zeta_5(s) &= 2k'\tau + k\tau', \\ \zeta_6(s) &= k^2\tau. \end{aligned}$$

*Proof.* From (3.1) we have

$$\nabla_T^2T = (-k^2)T + (k')N + (k\tau)B,$$

$$\nabla_T^3T = (-3kk')T + (k'' - k(k^2 + \tau^2))N + (2k'\tau + k\tau')B,$$

$$\nabla_T^5 T = \zeta_1(s)T + \zeta_2(s)N + \zeta_3(s)B.$$

By (3.10) we see that  $\psi$  is a III-harmonic curve if and only if

$$\zeta_1(s)T + \zeta_2(s)N + \zeta_3(s)B + \zeta_4(s)R(N, T)T + \zeta_5(s)R(B, T)T - \zeta_6(s)R(B, N)T = 0. \quad (3.12)$$

Using (2.2), we obtain (3.11).  $\square$

**Theorem 3.3.** *Let  $\psi : I \rightarrow \widetilde{\mathbf{SL}_2\mathbb{R}}$  be a differentiable curve parametrized by arc length. Then  $\psi$  is a proper non-geodesic III-harmonic curve if and only if the following equations hold,*

$$\begin{cases} \zeta_1(s) = 0 \\ \zeta_2(s) + \zeta_4(s)(1 - 8B_3^2) + 8\zeta_5(s)N_3B_3 + 8\zeta_6(s)T_3B_3 = 0 \\ \zeta_3(s) + 8\zeta_4(s)N_3B_3 + \zeta_5(s)(1 - 8N_3^2) - 8\zeta_6(s)T_3N_3 = 0. \end{cases} \quad (3.13)$$

*Proof.* Combining (3.4) and (3.11), we obtain (3.13).  $\square$

**Corollary 3.4.** *If  $\tau = 0$  and  $N_3B_3 \neq 0$ , then  $k = 0$ .*

*Proof.* From (3.13), we get

$$\begin{cases} \zeta_1(s) = 0 \\ \zeta_2(s) + (1 - 8B_3^2)(k'' - 2k^3) = 0 \\ N_3B_3(k'' - 2k^3) = 0. \end{cases}$$

Then the following hold:

$$2k'k'' + kk^{(3)} - 2k^3k' = 0, \quad (3.14)$$

$$k'' = 2k^3, \quad (3.15)$$

$$k^5 + k^{(4)} - 15kk'^2 - 10k^2k'' = 0. \quad (3.16)$$

The equation (3.14) is rewritten as

$$k'k'' + (kk'')' - 2k^3k' = 0. \quad (3.17)$$

From (3.15) and (3.17) we obtain  $k = \text{constant}$ . Hence (3.16) implies  $k = 0$ .  $\square$

## 4 III-harmonic helices in $\widetilde{\mathbf{SL}_2\mathbb{R}}$

A circular helix is a geometric curve with non-vanishing constant curvature and non-vanishing constant torsion. We give the following theorem

**Theorem 4.1.** *Let  $\psi : I \rightarrow \widetilde{\mathbf{SL}_2\mathbb{R}}$  be a non-geodesic proper III-harmonic curve with constant curvature  $k$  and nonzero torsion  $\tau$ . Then  $\psi(s)$  is a circular helix. Further,  $k$  and  $\tau$  satisfy the differential equation system:*

$$\begin{cases} (k^2 + \tau^2)^2 - (2k^2 + \tau^2)(1 - 8B_3^2) + 8k\tau T_3B_3 = 0 \\ N_3((2k^2 + \tau^2)B_3 + k\tau T_3) = 0. \end{cases} \quad (4.1)$$

*Proof.* From the first equation of (3.13), we obtain that  $\tau = \text{constant}$ . Then  $\psi(s)$  is a circular helix. If we substitute  $k = \text{constant}$  and  $\tau = \text{constant}$  in (3.13), we obtain (4.1).  $\square$

**Theorem 4.2.** *Let  $\psi : I \rightarrow \widetilde{\mathbf{SL}_2\mathbb{R}}$  be a non-geodesic III-harmonic helix parametrized by arc length. If  $N_3 = 0$ , then  $B_3 = \text{constant}$ .*

*Proof.* Using (3.5), we have

$$g_{\widetilde{\mathbf{SL}_2\mathbb{R}}}(\nabla_T B, e_3) = B'_3 - (T_1 B_2 - T_2 B_1).$$

Then from the Frenet equations (3.1) we obtain

$$g_{\widetilde{\mathbf{SL}_2\mathbb{R}}}(\nabla_T B, e_3) = -\tau N_3. \quad (4.2)$$

Since  $N_3 = 0$  and  $T_1 B_2 - T_2 B_1 = -N_3$ , we have

$$g_{\widetilde{\mathbf{SL}_2\mathbb{R}}}(\nabla_T B, e_3) = B'_3. \quad (4.3)$$

These, together with (4.2) and (4.3) give  $B'_3 = 0$ .  $\square$

**Corollary 4.3.** *Let  $\psi : I \rightarrow \widetilde{\mathbf{SL}_2\mathbb{R}}$  be a non-geodesic III-harmonic helix parametrized by arc length. If  $N_3 = 0$ , then  $T_3 = \text{constant}$ .*

*Proof.* By Theorem 4.2 and (4.1), we have  $T_3 = \text{constant}$ .  $\square$

**Theorem 4.4.** *Let  $\psi : I \rightarrow \widetilde{\mathbf{SL}_2\mathbb{R}}$  be a non-geodesic III-harmonic helix parametrized by arc length. If  $N_3 \neq 0$ , then*

$$\tau^2 = \frac{1 - 2k^2 - \sqrt{1 + 4k^2}}{2}, \quad \text{or}, \quad \tau^2 = \frac{1 - 2k^2 + \sqrt{1 + 4k^2}}{2}. \quad (4.4)$$

*Proof.* Since  $N_3 \neq 0$ , from (4.1), we obtain

$$(k^2 + \tau^2)^2 - (2k^2 + \tau^2) = 0. \quad (4.5)$$

Hence (4.5) implies (4.4).  $\square$

## 5 General helix in $\widetilde{\mathbf{SL}_2\mathbb{R}}$

In this section we give a characterization of general helices with respect to the Frenet frame  $\{T, N, B\}$  in  $\widetilde{\mathbf{SL}_2\mathbb{R}}$ . A necessary and sufficient condition for a curve to be a general helix is that the ratio of curvature to torsion be constant (see [1–3, 20, 24–26] for details).

**Definition 5.1.** *Let  $\psi(s)$  be a curve in  $\widetilde{\mathbf{SL}_2\mathbb{R}}$  and  $\{T, N, B\}$  be the Frenet frame on  $\widetilde{\mathbf{SL}_2\mathbb{R}}$  along  $\psi$ . A curve  $\psi$  such that*

$$\frac{\tau}{k} = \Lambda, \quad \Lambda \in \mathbb{R}, \quad (5.1)$$

*is called a general helix with respect to Frenet frame.*



**Remark 5.1.** i) If  $k = 0$ , then it indicates a straight line.

ii) If  $k = \text{constant} \neq 0$  and  $\tau = 0$ , then the curve  $\psi$  is a circle.

**Theorem 5.1.** Let  $\psi : I \rightarrow \widetilde{\mathbf{SL}_2\mathbb{R}}$  be a non-geodesic III-harmonic general helix parametrized by arc length. If  $N_3 = 0$ , then  $\psi$  is a circular helix and  $B_3 = \text{constant}$ .

*Proof.* From (5.1), we have

$$\begin{cases} \zeta_1(s) = -10k'k'' - 5kk^{(3)} + 10k^3k'(\Lambda^2 + 1) \\ \zeta_2(s) = (1 + \Lambda^2)^2k^5 + k^{(4)} - 5(1 + \Lambda^2)k(2kk'' + 3k'^2) \\ \zeta_3(s) = -\Lambda\xi_1(s) \\ \zeta_4(s) = k'' - k^3(\Lambda^2 + 2) \\ \zeta_5(s) = 3\Lambda k'k \\ \zeta_6(s) = \Lambda k^3. \end{cases} \quad (5.2)$$

By using equations (5.2) in (3.13), we can obtain a system of three differential equations characterizing III-harmonic general helix in  $\widetilde{\mathbf{SL}_2\mathbb{R}}$

$$\begin{cases} \zeta_1(s) = 0 \\ \zeta_2(s) + \zeta_4(s) + 8B_3(-\zeta_4(s)B_3 + \zeta_5(s)N_3 + \zeta_6(s)T_3) = 0 \\ N_3(\zeta_4(s)B_3 - N_3\zeta_5(s) - \zeta_6(s)T_3) + \zeta_5(s) = 0. \end{cases} \quad (5.3)$$

Substituting  $N_3 = 0$  into the third equation in (5.3) we have  $\zeta_5(s) = 0$ , which implies  $k = \text{constant}$  and hence  $\tau = \text{constant}$ . Then  $\psi$  is a circular helix.  $\square$

From (3.1) and (3.5), we have

$$g_{\widetilde{\mathbf{SL}_2\mathbb{R}}}(\nabla_TB, e_3) = -\tau N_3, \quad g_{\widetilde{\mathbf{SL}_2\mathbb{R}}}(\nabla_TB, e_3) = B'_3 + T_2B_1 - T_1B_2.$$

Since  $N_3 = T_2B_1 - T_1B_2 = 0$ , it follows that  $B_3 = \text{constant}$ .

From the second equation of (5.3), we have

**Corollary 5.2.** Let  $\psi : I \rightarrow \widetilde{\mathbf{SL}_2\mathbb{R}}$  be a non-geodesic III-harmonic general helix parametrized by arc length. If  $N_3 = 0$ , then  $T_3 = \text{constant}$ .

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Bendehiba Senoussi

Department of Mathematics, Higher Normal School of Mostaganem, Mostaganem, Algeria

E-mail: [se2014bendhiba@gmail.com](mailto:se2014bendhiba@gmail.com)

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